The Limiting Amplitude Principle for the Wave Equation with Variable Coefficients

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The Limiting Amplitude Principle (LAP): the classical and the present setups

$$\begin{cases} \partial_t^2 u(\mathbf{x}, t) + \mathcal{L}[\mathbf{u}](\mathbf{x}, t) = e^{-i\omega t} F(\mathbf{x}), & \mathbf{x} \in \mathbb{R}^d \setminus \overline{\Omega}_{int}, t > 0, \\ u(\mathbf{x}, t) = 0, & \mathbf{x} \in \partial \Omega_{int}, t > 0, \\ u(\mathbf{x}, 0) = 0, & \partial_t u(\mathbf{x}, 0) = 0, & \mathbf{x} \in \mathbb{R}^d \setminus \overline{\Omega}_{int}. \end{cases}$$
(in so

LAP (if valid):

$$u(x, t) \xrightarrow[t \to \infty]{} e^{-i\omega t} U(x)$$

n some sense and with some rate)





Frequency-domain problem

$$\begin{aligned} & \left[-\omega^2 U(\mathbf{x}) + \mathcal{L} \left[U \right](\mathbf{x}) = F(\mathbf{x}) , & \mathbf{x} \in \mathbb{R}^d \setminus \overline{\Omega}_{int}, \\ & U(\mathbf{x}) = 0, & \mathbf{x} \in \partial \Omega_{int}, \\ & \text{(Sommerfeld radiation condition),} & |\mathbf{x}| \to \infty. \end{aligned}$$

(4)

 $\blacktriangleright \text{ Classical setting: } d = 3, \quad \mathcal{L}[u] := -c_0^2 \Delta u, \quad \Omega_{int} \subset \mathbb{R}^3 \text{ open and bounded or } \Omega_{int} = \emptyset, \quad F \in L^2(\mathbb{R}^3 \setminus \overline{\Omega}_{int}) \text{ and compactly supported}, \quad c_0 > 0 \text{ constant}.$

 $\blacktriangleright \text{ Present setting: } d \in \{1, 2, 3\}, \quad \mathcal{L}[u] := -\beta^{-1}(\mathbf{x}) \nabla \cdot (\alpha(\mathbf{x}) \nabla u), \quad \Omega_{int} = \emptyset, \quad F \in L^2(\mathbb{R}^d) \text{ and compactly supported,}$ $\alpha, \beta \in C^{\infty}(\mathbb{R}^d), \quad \alpha(x), \beta(x) > 0 \text{ non-trapping} \text{ and } (\alpha(x) - \alpha_0), (\beta(x) - \beta_0) \text{ are compactly supported for some constants } \alpha_0, \beta_0 > 0.$

Sommerfeld radiation condition: $\lim_{|\mathbf{x}|\to\infty} |\mathbf{x}|^{\frac{d-1}{2}} \left| \partial_{|\mathbf{x}|} U(\mathbf{x}) - i\omega \sqrt{\beta_0/\alpha_0} U(\mathbf{x}) \right| = 0.$

Our results

Main features:

Time-domain problem

- The LAP is proved for the wave equation that has non-constant coefficients and which is not necessarily in a divergent form. Besides the classical case d = 3, we also consider d = 2.
- The validity of the LAP is extended to the case d = 1 with an appropriate modification.
- The convergence in the LAP is quantified and is shown to be algebraic in time for $d \in \{2, 3\}$ and exponential for d = 1.

Theorem (LAP: d = 2 and d = 3)

Let $d \in \{2,3\}$. Suppose that all assumptions of the present setting are satisfied. Let u(x, t) and U(x) be the solutions to (1) and (2), respectively. Then, for any bounded $\Omega \subset \mathbb{R}^d$, there exists a constant C > 0 such that

Some ingredients of the proof

Most of the auxiliary results concern time decay of the solutions of the following initial-value problems:

 $\int \partial_t^2 u(\mathbf{x}, t) - \beta^{-1}(\mathbf{x}) \nabla \cdot (\alpha(\mathbf{x}) \nabla u(\mathbf{x}, t)) = f(\mathbf{x}, t), \quad \mathbf{x} \in \mathbb{R}^d, \ t > 0,$ (3) $u(\mathbf{x},0) = u_0(\mathbf{x}), \quad \partial_t u(\mathbf{x},0) = u_1(\mathbf{x}), \quad \mathbf{x} \in \mathbb{R}^d,$ $\left\{ egin{aligned} &\partial_t^2 v\left(\mathrm{x},t
ight) - c_0^2 \Delta v\left(\mathrm{x},t
ight) = 0, \quad \mathrm{x} \in \mathbb{R}^d, \quad t > 0, \ & v\left(\mathrm{x},0
ight) = v_0\left(\mathrm{x}
ight), \quad \partial_t v\left(\mathrm{x},0
ight) = v_1\left(\mathrm{x}
ight), \quad \mathrm{x} \in \mathbb{R}^d, \end{aligned}
ight.$ where $c_0 := \sqrt{\alpha_0/\beta_0}$.

Lemma (for the solution of (3))

Let $d \geq 2$, $f \equiv 0$, $u_0 \in H^3(\mathbb{R}^d)$, $u_1 \in H^2(\mathbb{R}^d)$ and $\int_{\mathbb{R}^d} \left(1+|\mathbf{x}|^2\right)^{d+1+\epsilon} \left(|u_0(\mathbf{x})|^2+|u_1(\mathbf{x})|^2\right) d\mathbf{x} < \infty \text{ with some } \epsilon > 0.$ Then, for any bounded $\Omega \subset \mathbb{R}^d$, there exists a constant C > 0 such that $\|u(\cdot,t)\|_{H^1(\Omega)} + \|\partial_t u(\cdot,t)\|_{L^2(\Omega)} \leq \frac{C}{(1+t^2)^{\frac{d-1}{2}}}, \quad t \geq 0.$ Moreover, if d = 1, $f \equiv 0$, $u_0 \in H^1(\mathbb{R})$, $u_1 \in L^2(\mathbb{R})$ and are supported inside some $\Omega_0 \subset \mathbb{R}$. Then, for any bounded interval $\Omega \subset \mathbb{R}$, there exist a constant C > 0 such that

$$\begin{aligned} & \text{for } d = 2, \ t \ge 0: \\ & \left\| u\left(\cdot, t\right) - e^{-i\omega t} U \right\|_{H^{1}(\Omega)} + \left\| \partial_{t} u\left(\cdot, t\right) + i\omega e^{-i\omega t} U \right\|_{L^{2}(\Omega)} \le C \frac{1 + \log\left(1 + t^{2}\right)}{(1 + t^{2})^{\frac{1}{2}}} \end{aligned}$$
$$& \text{for } d = 3, \ t \ge 0: \\ & \left\| u\left(\cdot, t\right) - e^{-i\omega t} U \right\|_{H^{1}(\Omega)} + \left\| \partial_{t} u\left(\cdot, t\right) + i\omega e^{-i\omega t} U \right\|_{L^{2}(\Omega)} \le \frac{C}{(1 + t^{2})^{\frac{1}{2}}}.\end{aligned}$$

Theorem (LAP: d = 1 - an extension of the principle)

Let d = 1. Suppose that all assumptions of the present setting are satisfied. Let u(x, t) and U(x) be the solutions to (1) and (2), respectively. Then, for any bounded interval $\Omega \subset \mathbb{R}$, there exist constants $u_{\infty} \in \mathbb{C}$, $\Lambda > 0$ and C > 0 such that, for $t \geq 0$,

$$\left\| u\left(\cdot,t\right)-e^{-i\omega t}U-u_{\infty}\right\|_{H^{1}(\Omega)}+\left\| \partial_{t}u\left(\cdot,t\right)+i\omega e^{-i\omega t}U\right\|_{L^{2}(\Omega)}\leq Ce^{-\Lambda t}.$$

 $\|u(\cdot,t)-u_{\infty}\|_{H^{1}(\Omega)}+\|\partial_{t}u(\cdot,t)\|_{L^{2}(\Omega)}\leq Ce^{-\Lambda t}, \quad t\geq 0,$ where $u_{\infty} := \frac{1}{2} \int_{\Omega_0} \beta(x) u_1(x) dx$ and $\Lambda > 0$ can be estimated explicitly.

Lemma (for the solution of (3))

Let $d \in \{2,3\}$, $u_0 \equiv 0$, $u_1 \equiv 0$, $f \in C(\mathbb{R}_+, L^2(\mathbb{R}^d))$, $f(\cdot, t)$ is compactly supported for each $t \geq 0$, and, for some $C_0 > 0$, $\|f\left(\cdot,t
ight)\|_{L^2(\mathbb{R}^d)}+\|\partial_t f\left(\cdot,t
ight)\|_{L^2(\mathbb{R}^d)}\leq rac{\mathcal{L}_0}{\left(1+t^2
ight)^{rac{1}{2}}},\quad t\geq 0.$ Then, for any bounded $\Omega \subset \mathbb{R}^d$, there exists a constant C > 0 such that

A practical application

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One motivation to revisit the LAP and study rates of convergence is analysis of time-domain numerical methods.

In particular, such methods can be very attractive for high frequencies. For example, by an appropriate treatment, computational cost can be effectively reduced by directing numerical efforts towards the wavefront area yet also resolving the wavefield in the whole domain of interest:

Arnold, A., Geevers, S., Perugia, I., Ponomarev, D. An adaptive finite element method for high-frequency scattering problems with variable coefficients, arXiv:2103.02511, 2021.

 $\begin{array}{l} \text{for } d=2, \ t\geq 0: \ \|u\left(\cdot,t\right)\|_{H^{1}(\Omega)}+\|\partial_{t}u\left(\cdot,t\right)\|_{H^{1}(\Omega)}\leq C\frac{1+\log\left(1+t^{2}\right)}{\left(1+t^{2}\right)^{\frac{1}{2}}}.\\ \text{for } d=3, \ t\geq 0: \ \|u\left(\cdot,t\right)\|_{H^{1}(\Omega)}+\|\partial_{t}u\left(\cdot,t\right)\|_{H^{1}(\Omega)}\leq \frac{C}{\left(1+t^{2}\right)^{\frac{1}{2}}}.\end{array}$

Lemma (for the solution of (4))

Let $d \in \{2,3\}$ and v_0 , v_1 are of some special form: either sufficiently localised and oscillatory or essentially outgoing. Then, for any bounded $\Omega \subset \mathbb{R}^d$, there exists a constant C > 0 such that

$$|\mathbf{v}(\mathbf{x},t)| + \ldots + |\partial_t \Delta \mathbf{v}(\mathbf{x},t)| \leq rac{C}{(1+t^2)^{rac{1}{2}}}, \quad \mathbf{x} \in \Omega, \ t \geq 0$$

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